## ERRATA

## Erratum: Surface effects on spinodal decomposition in the framework of a linearized theory [Phys. Rev. E 52, 2848 (1995)]

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Due to an oversight, we neglected to note that our Eq. (2) in Ref. [1] should read

$$\frac{\partial}{\partial \tau} \delta \phi_{k\parallel}(0,\tau) = h_1 \delta(k_{\parallel},0) + \left[g + \overline{\sigma}_s k_{\parallel}^2\right] \delta \phi_{k\parallel}(0,\tau) + \gamma \frac{\partial}{\partial Z} \delta \phi_{k\parallel}(Z,\tau)|_{Z=0} - \left(\frac{\gamma}{4}\right)^{2/3} \frac{\partial^2}{\partial Z^2} \delta \phi_{k\parallel}(Z,\tau)|_{Z=0} - \frac{5}{6} \left(\frac{\gamma}{4}\right)^{1/3} \frac{\partial^3}{\partial Z^3} \delta \phi_{k\parallel}(Z,\tau)|_{Z=0}.$$
(2)





FIG. 4. Surface part V(Z,s) of the Laplace transform  $\overline{u}(Z,s)$  plotted vs the scaled distance Z for the case  $h_1=4$ ,  $\gamma=4$ , g=-4, amplitude  $u_0=0.025$ ,  $\phi_0=0$ ,  $\overline{\sigma_s}=4$ , and two values of the scaled wavenumber  $k_{\parallel}: k_{\parallel}=1$  (b) and  $k_{\parallel}=\sqrt{2}$  (c). In each case, four values of the scaled frequency *s* are shown, as indicated in the figure. Note the frequency limit  $s'_0=1$  here, while  $s_0=1$  (b) and 1/3 (c).

FIG. 5. Surface part V(Z,s) of the Laplace transform  $\overline{u}(Z,s)$  plotted vs the scaled distance Z for the case  $h_1=4$ ,  $\gamma=4$ , g=-4, amplitude  $u_0=0.025$ ,  $\phi_0=0.47$ ,  $\overline{\sigma_s}=4$ , and two values of the scaled wavenumber  $k_{\parallel}: k_{\parallel} = \sqrt{1-3\phi_0^2}$  (b) and  $k_{\parallel} = \sqrt{2}\sqrt{1-3\phi_0^2}$  (c). In each case, four values of the scaled frequency s are shown, as indicated in the figure.



FIG. 6. Same as Fig. 5 but for  $\phi_0 = 0.56$ .

In the original Eq. (2) the  $\delta(k_{\parallel},0)$  and the  $\overline{\sigma}_{s}k_{\parallel}^{2}\delta\phi_{k_{\parallel}}(0,\tau)$  were missing. These terms were correctly given in Ref. [2], where the coefficient  $\overline{\sigma}_{s}$  is specified. Earlier, such a gradient-square term can be found on the seventh line of Eq. (34) of Ref. [3].

Clearly, our solutions given in Ref. [1] are unaffected for  $k_{\parallel}=0$ ; formally, they are still the same if  $h_1$  is replaced by

 $h_1(k_{\parallel}) = h_1 \delta(k_{\parallel}, 0)$  and  $g(k_{\parallel}) = g + \overline{\sigma_s} k_{\parallel}^2$ . These changes only affect the amplitude functions A(s), B(s), and V(s); all other parts of the solution are unaffected. Only parts (b) and (c) of Figs. 4, 5, and 6 are affected by this change. These parts are replaced by the present ones.

Our results are in complete agreement with those quoted in Ref. [2].

[2] H.P. Fischer, P. Maass, and W. Dieterich, Phys. Rev. Lett. 79, 893 (1997).

<sup>[1]</sup> H.L. Frisch, P. Nielaba, and K. Binder, Phys. Rev. E 52, 2848 (1995).

<sup>[3]</sup> K. Binder and H.L. Frisch, Z. Phys. B 84, 403 (1991).